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## Identification of chaotic systems with hidden variables (modified Bock's algorithm)

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### 9 Abstract

10 We address the problem of estimating parameters of chaotic dynamical systems from a time series in a situation  
 11 when some of state variables are not observed and/or the data are very noisy. Using specially developed quantitative  
 12 criteria, we compare performance of the original multiple shooting approach (Bock's algorithm) and its modified ver-  
 13 sion. The latter is shown to be significantly superior for long chaotic time series. In particular, it allows to obtain accu-  
 14 rate estimates for much worse starting guesses for the estimated parameters.

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### 17 1. Introduction

18 The problem of mathematical modeling of complex systems from experimental observables is well-known in different  
 19 fields of science and practice and has multiple names such as “reconstruction of dynamical systems” in nonlinear science  
 20 [1] and “system identification” in statistics and control theory [2]. It has different aspects and can be formulated in dif-  
 21 ferent ways. Here, we consider the case when the structure of model equations is known a priori from “the first prin-  
 22 ciples”. It reads

$$\frac{dy}{dt} = \mathbf{f}(\mathbf{y}, \mathbf{c}), \quad (1)$$

25 where  $\mathbf{y}$  is  $D$ -dimensional state vector,  $\mathbf{c}$  is  $P$ -dimensional parameter vector. The task is *to estimate the unknown param-*  
 26 *eters*  $c_1, \dots, c_P$  from a time series—discrete sequence of values observed at subsequent time instants  $\{\eta_1, \dots, \eta_N\}$ , where  
 27 an observable  $\eta$  is assumed to be a function of state vector  $\mathbf{y}$  (possibly corrupted with measurement noise),  $N$  is a time  
 28 series length. Let us consider the case when  $\eta$  is a scalar, which is quite typical and the most complicated. Such a for-  
 29 mulation has been considered in a number of works not only for differential equations [3–6], but also for maps [7–15]. In  
 30 practice, it is encountered in chemical kinetics (rate constants estimation) [16], laser physics (rates of transition between  
 31 energy levels) [17], electric engineering (ferroelectric and semiconductor nonlinearities) [18,19], cell biology (description  
 32 of signaling pathways [20], neuron modelling [21]), etc.

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33 Construction of the so-called standard models [22] demands the time series of state vectors  $\mathbf{y}$  to be at hand, i.e., one  
 34 must reconstruct all  $D$  components from a scalar time series  $\{\eta_i\}$ . For example, the observable itself can serve as one of  
 35 model variables, while others may be obtained via differentiation or integration. However, for a model structure spec-  
 36 ified from the first principles, some of state variables cannot often be measured or reconstructed from observed data.  
 37 Such variables are usually called “hidden”. The presence of hidden variables makes reconstruction a much more com-  
 38 plex problem, because deficit of information about hidden variables (which have also to be included into the set of esti-  
 39 mated quantities) requires more sophisticated approaches for parameter estimation. Usually, maximal likelihood  
 40 principle is appealed to, but practically it reduces to a version of the least-squares method. In the case considered here,  
 41 the problem is formalised as follows. One searches for initial conditions  $\mathbf{s}$  and parameters  $\mathbf{c}$  which provide the smallest  
 42 least-squares difference between the appropriate components of a model orbit  $\mathbf{y}(t)$  and observed data  $\bar{\mathbf{y}}^l$ . The sum of  
 43 errors (2) involves only  $l$  non-hidden variables:

$$S(\mathbf{s}, \mathbf{c}) = \sum_{i=1}^N [\mathbf{y}^l(t_i, \mathbf{s}, \mathbf{c}) - \bar{\mathbf{y}}_i^l]^2 = \min, \quad (2)$$

47 where  $\bar{\mathbf{y}}_i^l$  are observed vectors,  $\mathbf{y}^l(t_i, \mathbf{s}, \mathbf{c})$  are  $l$ -dimensional vectors consisting of the corresponding model state variables.  
 48 Minimisation of (2) is performed with the aid of iterative algorithms for some “starting guesses” for  $\mathbf{s}$  and  $\mathbf{c}$ .

49 In the case of a chaotic time series, a model trajectory is very sensitive to initial conditions. Therefore, “relief” of the  
 50 cost function (2) is very complex for large  $N$  and exhibits a lot of local minima. Thus, the “attracting area” of global  
 51 minimum is very narrow, so that it is unlikely to find it with arbitrary starting guesses. In order to overcome this dif-  
 52 ficulty, a special method—multiple shooting approach (Bock’s algorithm)—was proposed [16,23]. Later, it was noticed  
 53 [24] that it also encounters significant difficulties and additional efforts are necessary to succeed, although systematic  
 54 investigation of this problem is still lacking. In this work, we develop special measures to quantify the performance  
 55 of different parameter estimation techniques. With their aid, we compare different versions of multiple shooting ap-  
 56 proach (Section 2). By considering noisy time series of exemplary chaotic systems, we demonstrate that a modified  
 57 Bock’s algorithm allowing discontinuity of a model trajectory is the most efficient.

58 Chaotic dynamics and deficit of a priori information about system parameter values are typical in practice. There-  
 59 fore, the task considered here is of significant practical interest. We note also that the methods analysed here give pos-  
 60 sibility not only to estimate parameters, but also to reconstruct the time courses of hidden variables, which cannot be  
 61 measured by other means. So, the identification (reconstruction, parameter estimation) procedure acts as a universal  
 62 indirect “measuring device”.

## 63 2. Parameter estimation methods for hidden variable case

### 64 2.1. Initial value approach

65 This “naive” method consists in minimisation of (2) directly, where  $N$  is the length of the entire original time series.  
 66 In practice,  $N$  should be large enough to allow one to extract necessary information from noisy data. Furthermore, the  
 67 time series should involve all relevant time scales of the modeled dynamics. But for a chaotic time series, exponential  
 68 sensitivity of model orbits to initial conditions  $\mathbf{s}$  makes the attracting area of the global minimum of (2) very narrow.  
 69 Therefore, the initial value approach encounters great difficulties. The disadvantages of this approach are clearly shown,  
 70 e.g., in [27,19], so we do not pay significant attention to it here.

### 71 2.2. “Multiple shooting based” approaches

72 The name takes its origin from an analogy with well-known numerical methods for solution of a boundary-value  
 73 problem in ordinary differential equations. Since the multiple shooting approach accepts a number of variations, we  
 74 call all of them “multiple shooting based” approaches while the original one [16] just Bock’s algorithm.

#### 75 2.2.1. Original Bock’s algorithm

76 It is a modification of initial value approach which allows an increase in the time series length  $N$  and the use of start-  
 77 ing guesses for parameters not so close to their true values. This is possible since the entire time series is divided into  $L$   
 78 segments ( $n$  is the length of a segment,  $N = Ln$ ) and initial conditions for each of them  $\mathbf{s}_1, \mathbf{s}_2, \dots, \mathbf{s}_L$  are considered as  
 79 additional arguments of  $S$  (as quantities to be estimated):

$$S(\mathbf{s}_1, \dots, \mathbf{s}_L, \mathbf{c}) = \sum_{i=1}^L \sum_{j=1}^n [\mathbf{y}^j(t_j, \mathbf{s}_i, \mathbf{c}) - \bar{\mathbf{y}}^j(i-1)n + j]^2 = \min, \quad (3)$$

83 Time moments  $\tau_i = t_{(i-1)n+1}$  correspond to  $\mathbf{s}_i$ . In order to avoid a great number of free estimated quantities, that in-  
 84 creases variances of the estimates, one imposes a constraint of model trajectory continuity over the entire observed  
 85 interval:

$$\mathbf{y}(\tau_{i+1}, \mathbf{s}_i, \mathbf{c}) = \mathbf{s}_{i+1}, \quad i = 1, \dots, L-1. \quad (4)$$

89 Minimisation of (3) under the constraints (4) is the problem of constrained multidimensional optimisation. For arbi-  
 90 trarily chosen starting guesses for parameters and initial conditions, the model trajectory consists of  $L$  “disconnected”  
 91 pieces. However, it becomes “more continuous” gradually, after each iteration of the minimisation procedure. Bock’s  
 92 algorithm coincides with the initial value approach for  $L = 1$  and  $n = N$ .

93 It was claimed [23] that Bock’s algorithm does not require “genuine” starting guesses. Meanwhile, experience shows  
 94 that *this is not typically the case*. The algorithm extends the modelling capabilities only in part, since the condition (4) is  
 95 very strong. Therefore, often only local minima of (3) can be found.

### 96 2.2.2. Segmentation technique

97 In order to estimate parameters more accurately from a longer time series (where original Bock’s algorithm cannot  
 98 be applied due to local minima problem), it is divided into  $l$  shorter segments and parameters are estimated from each  
 99 segment independently without any constraints. The estimates obtained are averaged:  $\bar{\mathbf{c}} = \frac{1}{l} \sum_{i=1}^l \mathbf{c}_i$ . Such an approach is  
 100 called “piecewise” or “segmentation” technique [24]. If Bock’s algorithm is used for each segment, it is reasonable to  
 101 call it “segmentation Bock’s algorithm”. The disadvantage of this method is that the parameter estimators may be  
 102 strongly biased (even asymptotically) since an estimate from each short segment may be biased, which is not eliminated  
 103 via averaging. Therefore, the segmentation technique gives low accuracy of estimates as compared to the original Bock’s  
 104 algorithm if the global minimum can be easily found for both methods.

### 105 2.2.3. Modified Bock’s algorithm

106 It is known from statistical theory, e.g., [28], that the use of the entire time series in maximum likelihood estimation  
 107 is preferable for obtaining unbiased estimators than segmentation approach. So, we suggest to pay attention to a mod-  
 108 ification of Bock’s algorithm that has been already applied in [17,25,26] for non-chaotic signals consisting of a number  
 109 of independent shot realisations as a technique for “multiple experiment approach” problem solution. It was also  
 110 briefly mentioned in [24]. The idea is to refuse the constraints (4) for several  $(v-1)$  time instants holding the same  
 111 parameter values  $\mathbf{c}$  for the entire time series. So, the initial conditions for the  $v$  time instants, including the first one,  
 112 become independent quantities to be estimated. We choose these instants equidistantly within the time series. Such  
 113 an approach involves two adjustable parameters: the number of segments  $v$  and the number of subsegments within each  
 114 segment  $L$  ( $N = vLn$ ). Subsegments are required to apply Bock’s algorithm within each of the  $v$  segments.

115 The modified approach is not widely applied so far, even though it should have a number of advantages. The fact  
 116 that a final model trajectory is discontinuous is not an indication that the model is “bad” but weakening of the con-  
 117 straints (4) may help to find global minimum and reasonable model when “strict” Bock’s algorithm is not feasible.

## 118 3. Comparative study in numerical experiment

### 119 3.1. Comparison technique

120 We compare the methods using gray-scale “convergence diagrams” on the planes of starting guesses for parameters  
 121  $c_{i_1}, c_{i_2}$  (Fig. 1). White points denote starting guesses for which the global minimum is achieved, i.e., quite accurate esti-  
 122 mates are obtained. Gray colour means starting guesses from which minimisation procedure converges to a number of  
 123 local minima, darker colour corresponds to stopping at local minima situated further from the true values. We normal-  
 124 ise starting guesses so that the centre of a diagram corresponds to genuine guesses, i.e., to the true values of parameters  
 125  $c_i^0$ . The normalised starting guesses are denoted  $b_i = (c_i - c_i^0)/c_i^0$ . The size of white area on the diagrams quantifies the  
 126 estimation method’s performance. The broader this area, the better the method. Such areas typically have a very com-  
 127 plex structure (e.g., Fig. 1a), therefore we suggest an integral measure which is relative number  $\mu$  of white points within  
 128 a circle of radius  $r$ . The larger  $\mu$  (for a given  $r$ ), the better the method. We denote  $r_\mu$  the maximum value of the circle  
 129 radius corresponding to the relative ratio of white points equal to  $\mu$ . Here, we use mainly the value of  $r_{100}$ , which is the  
 130 radius of “100% convergence” to global minimum.

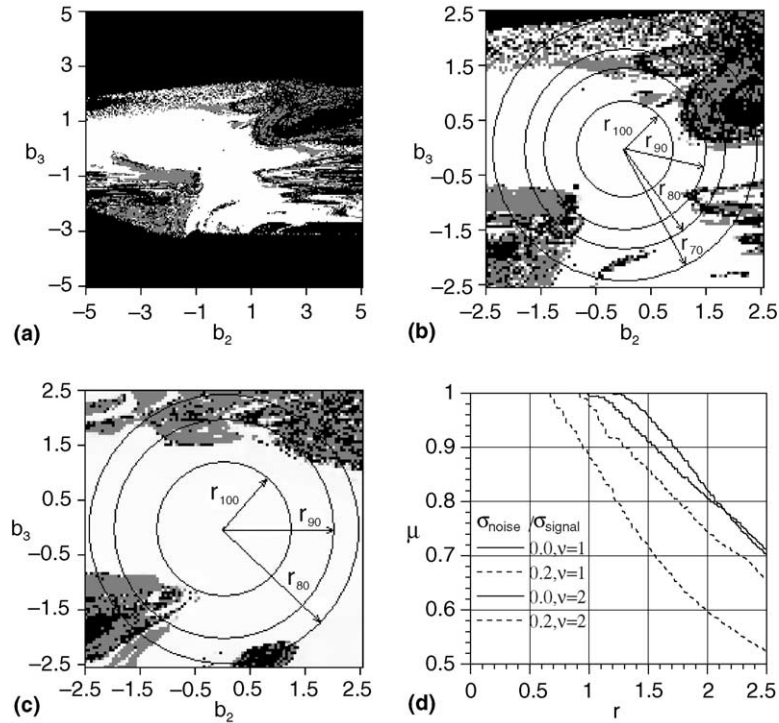


Fig. 1. The plane of normalised starting guesses for parameters of the Lorenz system (section with the plane  $b_1 = 0$ ). (a) Bock's algorithm with  $L = 30$ ,  $n = 35$ ; (b) is a magnification of (a); (c) the modified method with  $L = 15$ ,  $n = 35$ ,  $v = 2$ ; (d) the dependence  $\mu(r)$  for Bock's algorithm (black) and the modified method (gray) at different noise levels.

131 Below, we consider the case of three unknown parameters. So, three-dimensional diagrams for all three starting  
 132 guesses for parameters would contain complete information about the method's performance. Nevertheless, we use  
 133 two-dimensional projections for simplicity of illustration taking into account that they lead to the same qualitative con-  
 134 clusions about the methods' inferiority/superiority.

### 135 3.2. Identification of the Lorenz system

136 As the first test system for investigation of the performance of different parameter estimation techniques in case of  
 137 long chaotic time series and different starting guesses, we choose the Lorenz system

$$\dot{y}_1 = c_1(y_2 - y_1), \quad \dot{y}_2 = -y_2 + y_1(c_3 - y_3), \quad \dot{y}_3 = -c_2 y_3 + y_1 y_2, \quad (5)$$

140 with parameters  $c_1 = 10$ ,  $c_2 = 8/3$ ,  $c_3 = 46$  corresponding to a chaotic regime, and initial conditions  $y_1 = -7.60$ ,  
 141  $y_2 = -12.37$ ,  $y_3 = 38.66$  chosen arbitrarily on the chaotic attractor. The largest Lyapunov exponent is equal here to  
 142  $\lambda_1 = 1.23$  [23]. The equations are integrated with the fourth-order Runge–Kutta technique with stepsize 0.001 and sam-  
 143 pling interval 0.002 to generate a time series. An observed scalar time series is a realisation of the variable  $y_1$  corrupted  
 144 with additive Gaussian white noise:  $\eta = y_1 + \zeta$ . The variables  $y_2$  and  $y_3$  are regarded hidden.

145 Since the choice of genuine starting guesses for the values of  $y_2$  and  $y_3$  is unrealistic, we use the observable values as  
 146 starting guesses for all state variables  $s_1, s_2, \dots, s_L, \dots, s_{vL}$ . Even though such a choice is not the best possible, it is simple  
 147 and sometimes efficient [23]. To minimise the function (3) the generalised Gauss–Newton method is used [23].

148 Convergence of the original Bock's algorithm and the modified method to global minimum is illustrated in Fig. 1a  
 149 and b. These results correspond to the time series length for which the Bock's approach exhibits the best performance  
 150 (the broadest convergence region). Only the section of starting guesses space with the plane  $b_1 = 0$  is shown since un-  
 151 lucky choice of  $b_1$  is not so crucial as the choice of  $b_2, b_3$ . It can be seen that the area of 100% convergence of Bock's  
 152 algorithm is broad and the radius  $r_{100}$  is greater than 1.0, so relative deviations of starting guesses from true values (let  
 153 us call them errors in starting guesses) may exceed 100%. There is also a wide area which is very distant from global  
 154 minimum but allows to find global minimum (Fig. 1a). However, the modified method allows larger errors in starting

155 guesses as it can be seen from comparison of Fig. 1b and c. The values of  $r_{100}$ ,  $r_{90}$ , and  $r_{80}$  are greater for the modified  
 156 method and the white area is wider.

157 The value of  $\mu(r)$  for different noise levels is shown in Fig. 1d. The performance of both methods remains almost  
 158 unchanged for moderate noise. The horizontal line of 100% convergence ( $\mu = 1$ ) becomes shorter but not significantly:  
 159 in a noise-free setting its length is 1.2 for the modified approach and 1.1 for Bock's method, while for 20% noise-to-  
 160 signal ratio (ratio of rms amplitudes) it is 0.9 and 0.7, respectively. Similar conclusions can be drawn from Fig. 3c,  
 161 where the dependencies  $r_{100}(N)$  (for  $\nu = 1$  and  $\nu = 4$ ) are shown with black for noise-free setting and with gray for noisy  
 162 case.

163 The dependence  $r_{100}(L, n)$  shown in Fig. 2 also demonstrates the advantage of the modified method. Darker colour  
 164 corresponds to smaller values of  $r_{100}$  (they are indicated on the contour lines) at given starting guesses. For the modified  
 165 method, not only the area with  $r_{100} \geq 1$  is larger, but also there is an area where  $r_{100} \geq 1.2$  inside of it. This advantage  
 166 takes place for longer time series that is revealed by white hyperboles  $N = \text{constant}$  which are the lines of constant time  
 167 series length.

168 This conclusion is confirmed by Fig. 3a where the 100% convergence radius is shown versus time series length  $N$  for  
 169 different number of segments  $\nu$ . The number of subsegments  $L$  has been selected to make  $r_{100}$  as large as possible by  
 170 the use of hyperboles (Fig. 2) and choice of points from lighter areas. Hill-like shape of plots  $r_{100}(N)$  is determined by two  
 171 factors. For small  $N$ , the amount of data is insufficient to "average out" the noise influence, while for large  $N$ , the expo-  
 172 nential sensitivity to initial conditions takes place (small initial perturbations reaches the magnitude comparable to the  
 173 size of the attractor during time interval  $\tau_A = 1/\lambda_1$ ) that leads to complication of the cost function "relief". The curves  
 174 for larger  $\nu$  attain larger values of  $r_{100}$ , i.e., the modified method is more efficient than the original Bock's algorithm.  
 175 Those curves correspond also to larger values of  $N$ , therefore they are located closer to the right-hand side of the panel.  
 176 Furthermore, the range of time lengths within which the modified method is "100% convergent" increases with the  
 177 number of discontinuity points  $\nu$ , so the curves for greater  $\nu$  are "wider".

178 The investigation reveals (Fig. 3b) that the optimal value of segment length  $Ln$  is connected with Lyapunov time  $\tau_A$ .  
 179 Optimal time series lengths correspond to 1–2 Lyapunov times, see the upper horizontal axis in Fig. 3b. It is explained  
 180 as follows. The success of estimation depends on the segment length  $Ln$  (over which small initial perturbations of the  
 181 model orbit should not increase too strongly, so  $Ln$  should not be very large) and also on the number  $P + \nu D$  of free  
 182 parameters to be estimated (this number should not be very large since in very high-dimensional space relief of the cost  
 183 function may become very complicated also, i.e.,  $Ln$  should not be very small). As a consequence, there exists some  
 184 intermediate optimal value of  $Ln$  related via a certain proportionality constant to the characteristic time scale  $\tau_A$  of  
 185 the divergence of nearby model trajectories.

186 Fig. 3d shows the dependence of  $r_{100}$  on  $L$ , given a certain  $Ln$ . At that, there is also an optimal value of  $L$  as usually  
 187 for Bock's algorithm within each segment. The greatest  $r_{100}$  is achieved here for  $\nu = 2$  since greater  $\nu$  correspond just to  
 188 longer time series.

189 Similar results have been obtained from time series generated at different initial conditions, from time series of the  
 190 variable  $y_2$ , and from time series of  $y_1$  generated at a different set of "true" parameter values  $c_1 = 10$ ,  $c_2 = 8/3$ ,  $c_3 = 28$   
 191 that is known as a "classical" chaotical set for the Lorenz system.

192 We also had studied the jumps allowed by modified approach in points of discontinuity and we showed that these  
 193 jumps are small in comparison with attractor size: they are about  $10^{-3}$  from signal standard deviation even if 1% noise

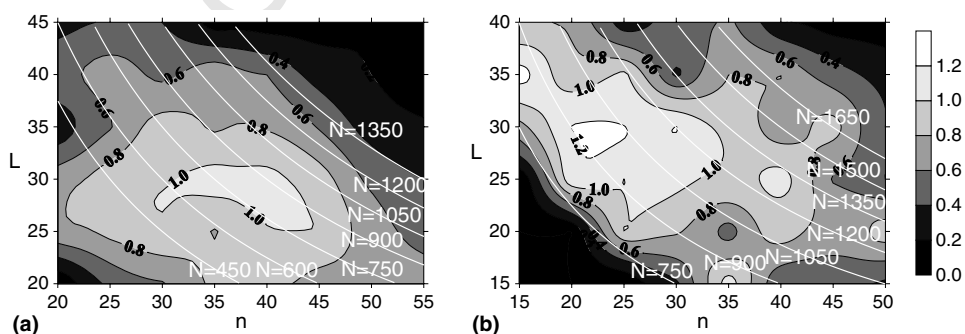


Fig. 2. The dependence  $r_{100}(L, n)$ : (a) for Bock's algorithm, (b) for the modified method with  $\nu = 2$ . Darker areas correspond to less radius  $r_{100}$ . The values of  $r_{100}$  are shown on the border lines. The white hyperboles are the lines of constant time series length  $N = \text{constant}$ .

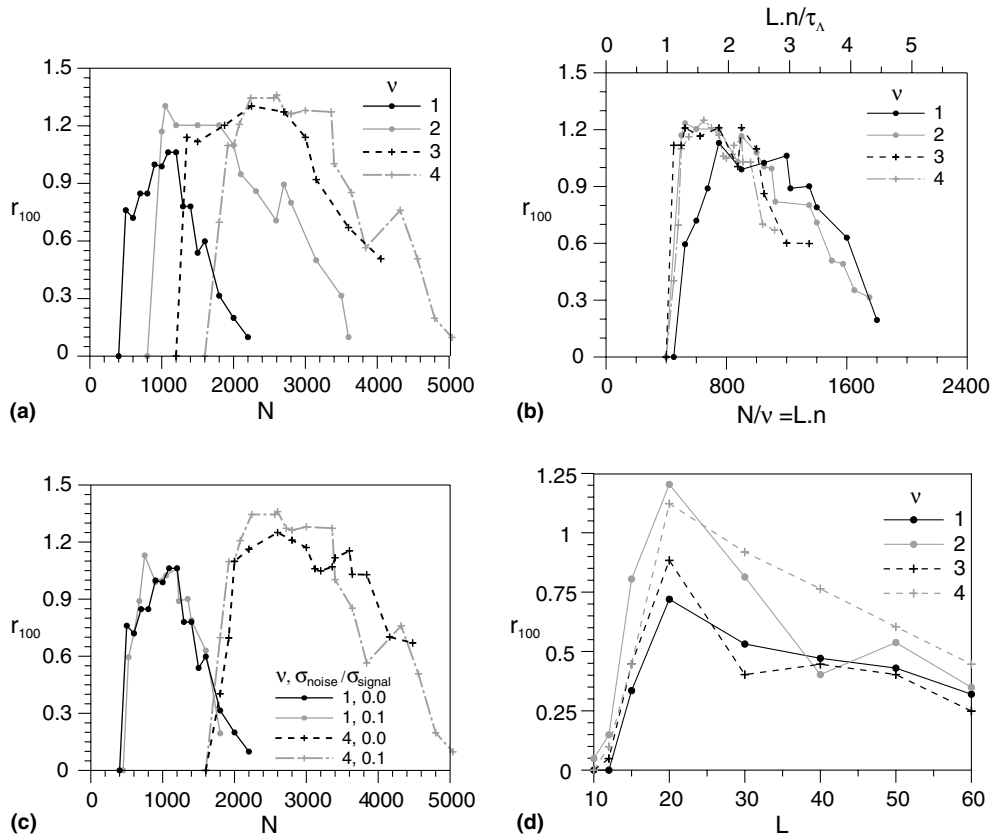


Fig. 3. The dependence of 100% convergence radius  $r_{100}$  on different factors: (a) on the entire time series length  $N$  for different numbers of discontinuity points allowed; (b) on the continuity segment length  $Ln$ ; (c) on the entire time series length  $N$  for different noise levels. (d) The dependence  $r_{100}(L)$  for different  $v$  and fixed continuity segment length  $Ln = 600$ .

194 is added to the observable. Though the jumps due to original Bock's approach are greatly smaller:  $10^{-10}$ – $10^{-12}$  from  
 195 signal standard deviation. It also has to be noticed that these "original" jumps decrease than discontinuity is allowed  
 196 so it can be said that the model imperfection is concentrated from the whole set of nodes to the nodes between the  
 197 segments.

### 198 3.3. Identification of Rössler system

199 In order to check whether our results hold for other systems, we perform the same investigation for the Rössler's  
 200 system.

$$\dot{y}_1 = -y_2 - y_3, \quad \dot{y}_2 = y_1 + c_1 y_2, \quad \dot{y}_3 = c_2 + y_3(y_1 - c_3), \quad (6)$$

204 with parameters  $c_1 = 0.2$ ,  $c_2 = 0.15$ ,  $c_3 = 10$ , that corresponds to a chaotic regime and initial conditions  $y_1 = 0.21$ ,  
 205  $y_2 = 6.5$ ,  $y_3 = 0.022$ . The basic "period" of oscillations is 6.0, the largest Lyapunov exponent is  $\lambda_1 = 0.1$ . The equations  
 206 (6) are integrated with fourth-order Runge–Kutta technique with stepsize 0.0002 and sampling interval 0.01. The vari-  
 207 able  $y_1$  is used as an observable both in a noise free setting and corrupted with additive Gaussian white noise.

208 We have chosen this system as an object since the "shape" of its attractor differs from the Lorenz one. The Lorenz  
 209 system oscillates near one of the two unstable fixed points in turn with irregular switchings between them. The simul-  
 210 taneous values of its  $y_1$  and  $y_2$  variables are relatively close to each other. Their "shift by a quarter of rotation period" is  
 211 a relatively small effect in absolute value as compared to the switchings between the two wings. The dynamics on the  
 212 Rössler attractor is a rotation about a single unstable fixed point (in projection onto the plane  $y_3 = 0$ ). So that the vari-  
 213 ables  $y_1$  and  $y_2$  are shifted in time by a quarter of the rotation period which is the main time scale here.

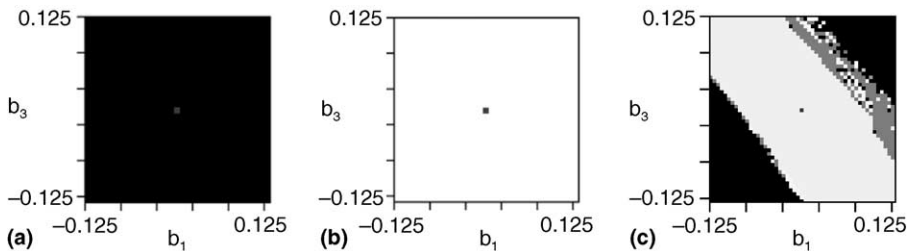


Fig. 4. The plane of normalised starting guesses for parameters of the Rössler system (section with the plane  $b_2 = 0$ ) illustrating convergence of the original Bock's algorithm: (a) all starting guesses for the hidden variables are equal to simultaneous observable values; (b) genuine starting guesses; (c) starting guesses are obtained via the time shift of the observed time series by a quarter of basic period.

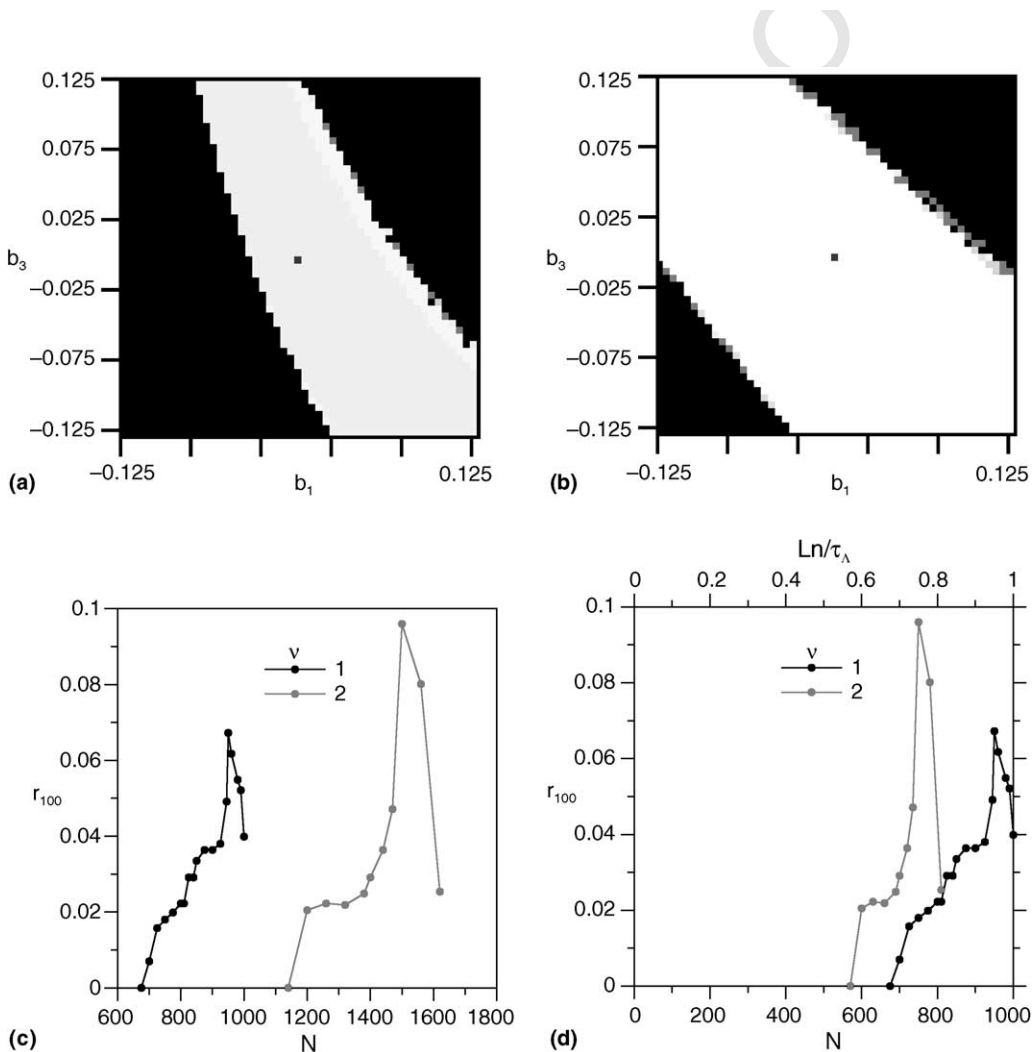


Fig. 5. The plane of normalised starting guesses for parameters of the Rössler system (section with the plane  $b_2 = 0$ ): (a) Bock's algorithm; (b) the modified method. (c) The dependance of  $r_{100}$  on the entire time series length  $N$  for the best choice of  $L$ . (d) The dependance of  $r_{100}$  on the segment length  $Ln$  for the best choice of  $L$ .

214 Due to such relationships between the state variables, the choice of starting guesses for the hidden variables equal to  
 215 the simultaneous observable value is more or less appropriate for the Lorenz system (as we have shown above) but leads

216 to unsuccessful results of parameter estimation in the Rössler system using any of the estimation techniques considered.  
217 In Fig. 4a it is shown that  $r_{100} = 0$ , i.e., one cannot find the global minimum for such a choice of starting guesses for the  
218 hidden variables at all. Quite good results are achieved if one uses genuine starting guesses for the hidden variables  
219 (Fig. 4b). To develop “good” and realistic starting guesses is also possible if one takes into account the knowledge  
220 about character of the original dynamics which can be gained by studying model dynamics. Namely, for the Rössler  
221 system it is relevant to take the observed time series shifted by a quarter of basic period as a starting guess for the var-  
222 iable  $y_2$  and zero as a starting guess for  $y_3$  because due to attractor features this variable is close to zero most of the time  
223 (Fig. 4c).

224 For starting guesses we proposed, the results of investigation are similar to that presented above for the Lorenz sys-  
225 tem and are shown in Fig. 5. They indicate that the modified method is successful in finding global minimum given  
226 starting guesses for parameters very far from the true values (Fig. 5b) while the original Bock’s algorithm demands more  
227 lucky starting guesses (Fig. 5a). Fig. 5c shows the dependence on the time series length  $N$  analogously to Figs. 3a and 5d  
228 shows the dependence on the segment length  $Ln$  analogously to Fig. 3b. The curves corresponding to larger  $\nu$  are  
229 “wider” and shifted to the right, i.e., the range of time series length allowing accurate estimation is greater for them.  
230 This advantage is observed for relatively long series that is similar to the results obtained for the Lorenz system.

#### 231 4. Conclusions

232 We compared performance of different methods for estimation of parameters (identification) of dynamical systems  
233 from chaotic time series in the case of hidden variables. All the methods rely upon the multiple shooting idea. The com-  
234 parison is done by using specially developed quantitative measure and considering exemplary chaotic systems. The ori-  
235 ginal Bock’s algorithm is shown to be less efficient than its modified version, which allows a model orbit to be  
236 discontinuous in several points within an observation interval.

237 The length of a time series and the number of its segments are shown to have significant influence upon the estima-  
238 tion results, and the choice of starting guesses for the hidden variables is quite important too. The chances for accurate  
239 estimation rise with time series length if the number of allowable points of model trajectory discontinuity is also in-  
240 creased. The optimal length of a continuity segment is close to Lyapunov time for long chaotic time series.

241 The modified method has a number of advantages as compared to the original Bock’s algorithm since it is not so  
242 demanding with respect to starting guesses for the hidden variables. This is due to weakening the model orbit continuity  
243 constraint. Therefore, longer time series can be processed with the modified method that allows to increase accuracy of  
244 the estimates. Moreover, the modified method is not so demanding with respect to starting guesses for parameters also,  
245 sometimes providing an opportunity to get accurate estimates when the original Bock’s algorithm fails.

246 The effect of measurement noise is shown to be not dramatical for both methods, even if noise-to-signal ratio is as  
247 high as 20% in rms amplitude. Note that if the time series length is fixed and the global minimum can be easily found for  
248 any estimation method then the accuracy of parameter estimates is the best for the original Bock’s algorithm, a bit  
249 worse for the modified method, and the worst for the segmentation technique. However, since global minimum can  
250 never be equally easily found for any method, the modified method should be considered as the best one from practical  
251 point of view. Finally, one should be careful when using the modified method, since specifying too small length of a  
252 model orbit continuity segments may lead to the situation where even a model with “incorrect”, “alien” structure is  
253 successfully fitted to observed data and erroneous conclusion about model adequacy is drawn.

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